A Survey on SPH Methods in Computer Graphics

MATERIALS

Dan Koschier Jan Bender Barbara Solenthaler Matthias Teschner



- Motivation
- Viscous Force
- Explicit Viscosity
- Implicit Viscosity
- Results

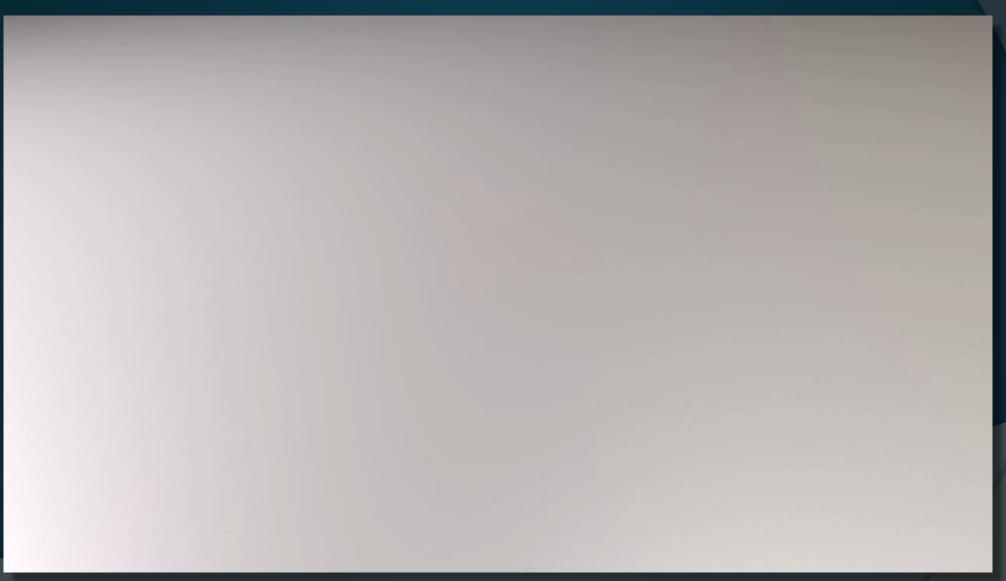
Motivation

Moored buoys

chain contacts are fully simulated with our proposed rigid body solver using only rigid-rigid contacts

> up to 5M fluid particles 2M static rigid particles 149k dynamic rigid particles

Viscous Fluids



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Classical Newtonian Fluid Model

Linear momentum equation:

$$\rho \frac{D\mathbf{v}}{Dt} = \nabla \cdot \mathbf{T} + \mathbf{f}$$

- > Density ρ
- \succ Velocity ${f v}$
- \succ Stress tensor $\, {f T} \,$
- \succ Force density f

Newtonian constitutive model:

$$\mathbf{T} = -p\mathbb{1} + 2\mu\mathbf{E}$$
$$\mathbf{E} = \frac{1}{2}\left(\nabla\mathbf{v} + (\nabla\mathbf{v})^T\right)$$

- \succ Strain rate tensor ${f E}$
- \succ Pressure p
- \succ Identity $\mathbb 1$
- > Dynamic viscosity μ

Navier-Stokes equations:

$$\rho \frac{D\mathbf{v}}{Dt} = -\nabla p + \mu \nabla^2 \mathbf{v} + \mathbf{f}$$

Viscous Force

• Viscous force for incompressible fluids:

$$\mathbf{f}_{\mathrm{visco}} = \mu \nabla^2 \mathbf{v}$$

 Recent SPH solvers either compute the divergence of the strain rate E or directly determine the Laplacian of v.

 Note that strain rate based approaches must enforce a divergencefree velocity field to avoid undesired bulk viscosity.

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Explicit Viscosity

• Standard SPH discretization of this Laplacian:

$$\nabla^2 \mathbf{v}_i = \sum_j \frac{m_j}{\rho_j} \mathbf{v}_j \nabla^2 W_{ij}$$

- Disadvantages:
 - The Laplacian of the kernel changes its sign inside the support radius
 - Sensitive to particle disorder [Mon05, Pri12]

Explicit Viscosity

 Alternative: determine one derivative using SPH and the second one using finite differences.

$$\nabla^2 v_i \approx 2(d+2) \sum_j \frac{m_j}{\rho_j} \frac{\mathbf{v}_{ij} \cdot \mathbf{x}_{ij}}{\|\mathbf{x}_{ij}\|^2 + 0.01h^2} \nabla W_{ij}$$
$$\mathbf{x}_{ij} = \mathbf{x}_i - \mathbf{x}_j, \ \mathbf{v}_{ij} = \mathbf{v}_i - \mathbf{v}_j, \ d = \text{dimension}$$

- Advantages:
 - Galilean invariant
 - vanishes for rigid body rotation
 - conserves linear and angular momentum



 Core idea of XSPH: reduce the particle disorder by smoothing the velocity field:

$$\hat{\mathbf{v}}_i = \mathbf{v}_i + \alpha \sum_j \frac{m_j}{\rho_j} (\mathbf{v}_j - \mathbf{v}_i) W_{ij}$$

- XSPH can also be used as artificial viscosity model.
- Advantage: The second derivative is not needed.
- Disadvantage: α is not physically meaningful.

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Takahashi et al. 2015

• Compute the viscous force as divergence of the strain rate $\,{f E}$:

$$\mathbf{f}_{\text{visco}} = \mu \nabla \cdot \left(\nabla \mathbf{v}_i + (\nabla \mathbf{v}_i)^T \right) = \mu \sum_j m_j \left(\frac{2\mathbf{E}_i}{\rho_i^2} + \frac{2\mathbf{E}_j}{\rho_j^2} \right) \nabla W_{ij}$$

• Implicit integration scheme

$$\mathbf{v}(t + \Delta t) = \mathbf{v}^* + \frac{\Delta t}{\rho} \mathbf{f}_{\text{visco}}(t + \Delta t)$$

• Solve linear system using the conjugate gradient method (CG).

Peer et al. 2015/2016

- Decompose velocity gradient: $\nabla \mathbf{v} = \mathbf{R} + \mathbf{V} + \mathbf{S}$
- Reduce shear rate by user-defined factor $0 \le \xi \le 1$:

 $\nabla \mathbf{v}^{\text{target}} = \mathbf{R} + \mathbf{V} + \xi \mathbf{S}$

• Reconstruct velocity field by solving linear system with CG:

$$\mathbf{v}_{i}(t + \Delta t) = \frac{1}{\rho_{i}} \sum_{j} m_{j} \left(\mathbf{v}_{j}(t + \Delta t) + \frac{\nabla \mathbf{v}_{i}^{\text{target}} + \nabla \mathbf{v}_{j}^{\text{target}}}{2} \mathbf{x}_{ij} \right) W_{ij}$$

Bender & Koschier 2017

• Define velocity constraint for each particle with user-defined factor:

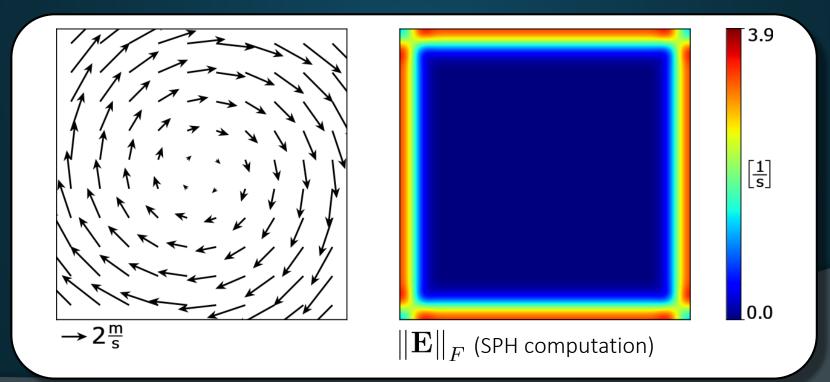
$$\mathbf{C}_i(\mathbf{v}) = \mathbf{E}_i - \gamma \mathbf{E}_i = \mathbf{0}, \quad 0 \le \gamma \le 1$$

• The constraint is a 6D function due to the symmetric strain tensor.

• Solve linear system for corresponding Lagrange multipliers.

Strain Rate Computation

- The introduced methods are based on the strain rate.
- However, computing the strain rate using SPH leads to errors at the free surface due to particle deficiency.



Weiler et al. 2018

• Implicit integration scheme

$$\mathbf{v}(t + \Delta t) = \mathbf{v}^* + \frac{\Delta t}{\rho} \mu \nabla^2 \mathbf{v}(t + \Delta t)$$

• Compute Laplacian as

$$\nabla^2 \mathbf{v}_i = 2(d+2) \sum_j \frac{m_j}{\rho_j} \frac{\mathbf{v}_{ij} \cdot \mathbf{x}_{ij}}{\|\mathbf{x}_{ij}\|^2 + 0.01h^2} \nabla W_{ij}$$

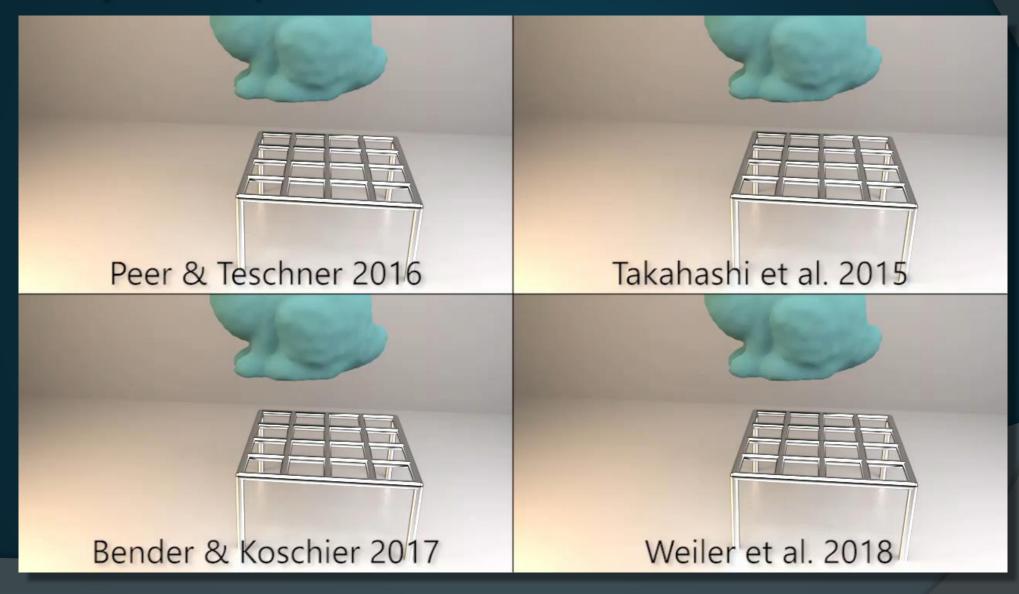
Solve linear system using a meshless conjugate gradient method.Laplacian approximation avoids problems at the free surface.

Discussion

- The strain rate based formulations lead to errors and artifacts at the free surface, which is avoided by directly computing the Laplacian of the velocity field.
- The viscosity parameters of Bender and Peer depend on the temporal and spatial resolution.
- Peer's reconstruction of the velocity field is fast but introduces a significant damping => not suitable for low viscous flow
- Takahashi et al. require the second-ring neighbors => low performance

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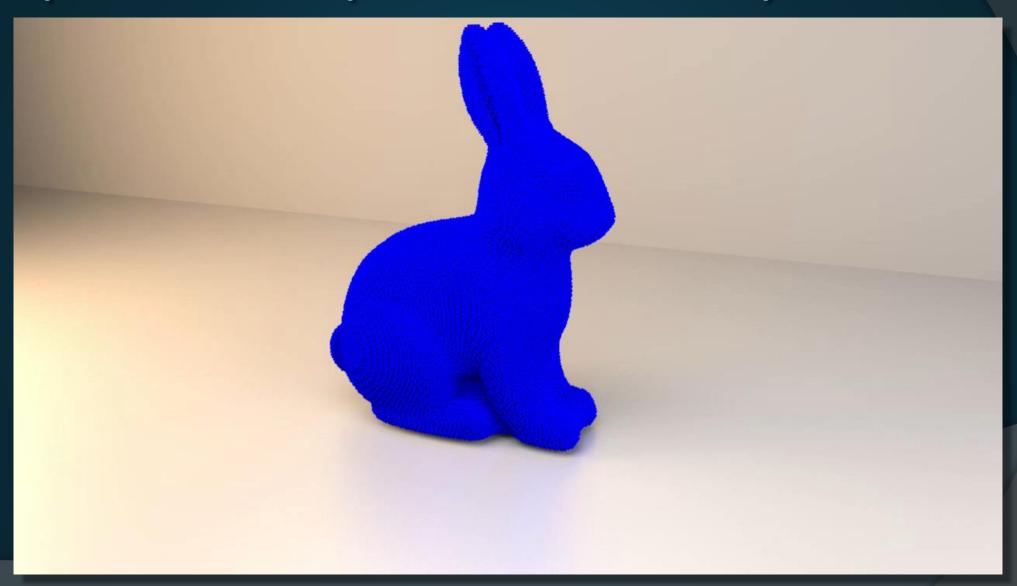
Quality Comparison



Coiling

Peer & Teschner 2016 Bender & Koschier 2017 Takahashi et al. 2015 Weiler et al. 2018

Temperature-Dependent Viscosity

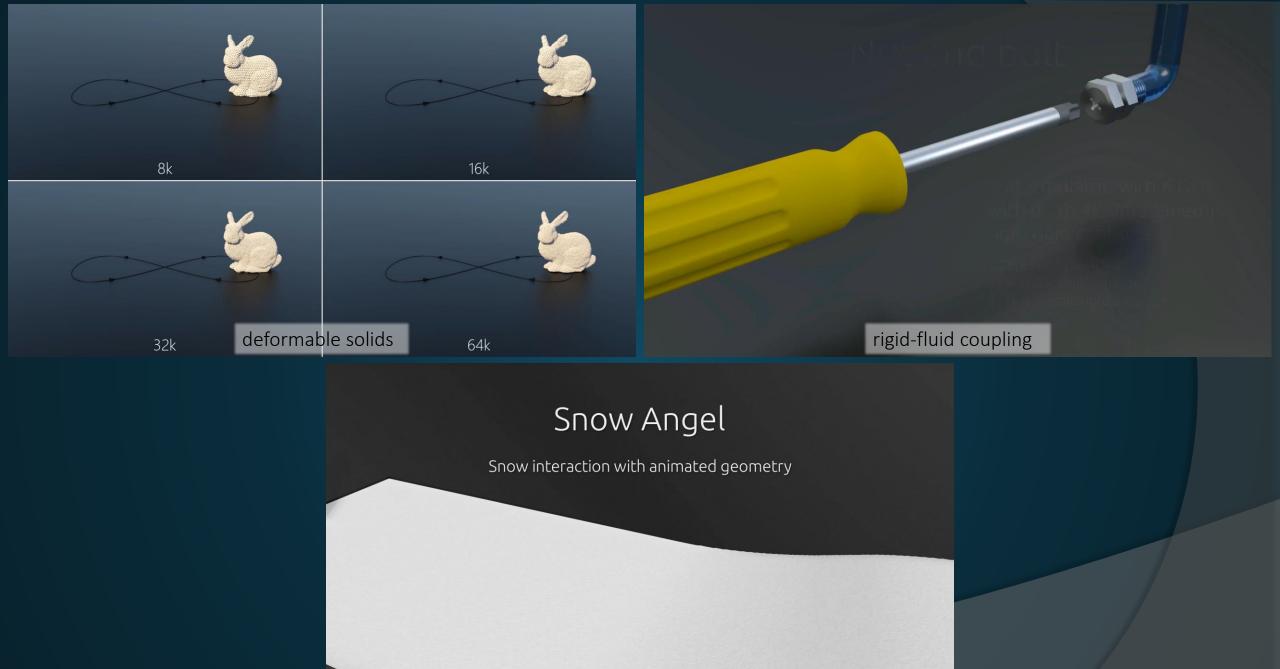


Summary

- Low viscous flow
 - Explicit methods are cheap and well-suited
 - Approximation of Laplacian yields better results while XSPH is slightly faster

• Highly viscous fluids

- Implicit methods are recommended to guarantee stability
- Strain rate based SPH formulations lead to artifacts at the free surface
- A formulation based directly on the Laplacian avoids this problem and generate more realistic results.



Want to try it out?

• SPlisHSPlasH is our open-source SPH library (MIT License)

- explicit and implicit viscosity methods
- various pressure solvers
- different boundary handling approaches
- effects: surface tension, vorticity, air drag
- multi-phase simulations
- deformable solids
- static and dynamic rigid bodies
- emitters

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- Python scripting
- exporters to Maya, Blender, Paraview



https://github.com/InteractiveComputerGraphics/SPlisHSPlasH