

A Survey on SPH Methods in Computer Graphics

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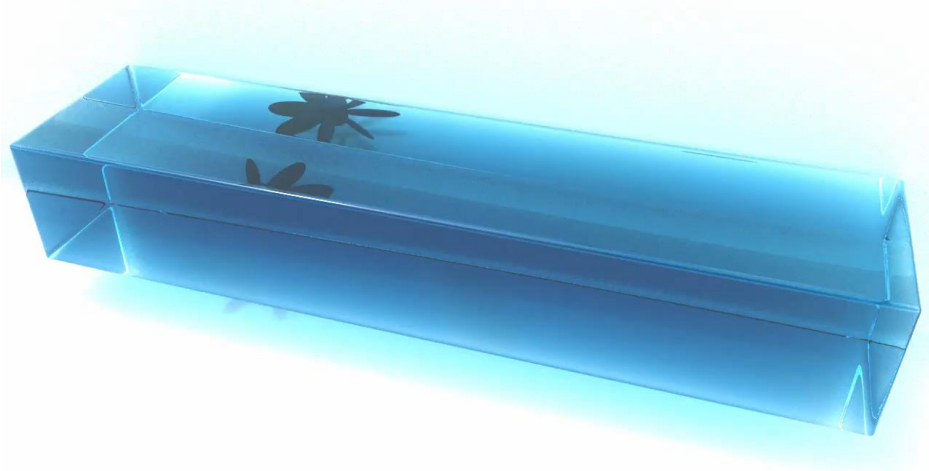


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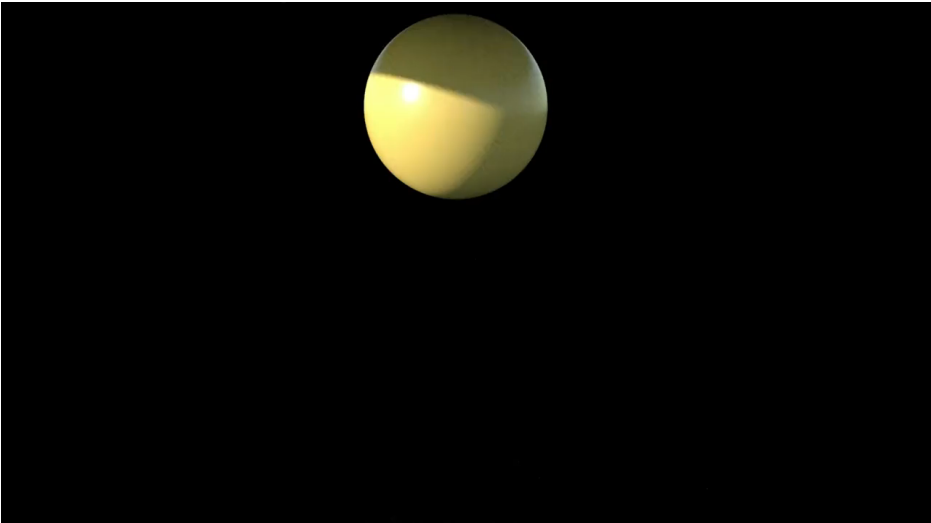
State-of-the-Art



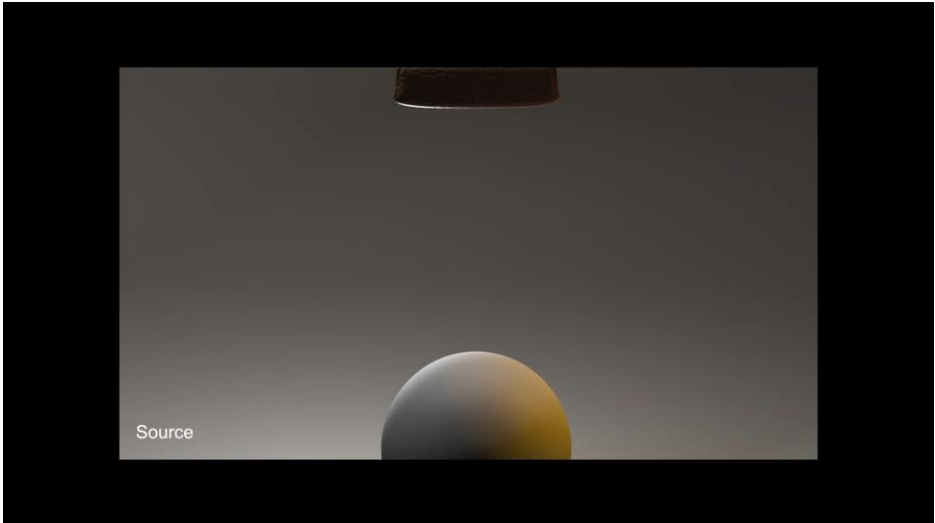
Bender et al. 2017



Weiler et al. 2018



Peer et al. 2018



Kim et al. 2020

Overview

Introduction



**Foundations of SPH and
PDE Discretization**

Material Models



**Incompressibility and
Pressure Solvers**

**Data-Driven Approaches
and ML Models**



Conclusion

Myths about SPH

“An SPH particle represents a molecule”

“SPH suffers from energy loss due to the smoothing”

“SPH is better than grid approaches”

“Grid approaches are better than SPH”

“SPH is (only) 0th-order consistent”

“An SPH particle represents a droplet”



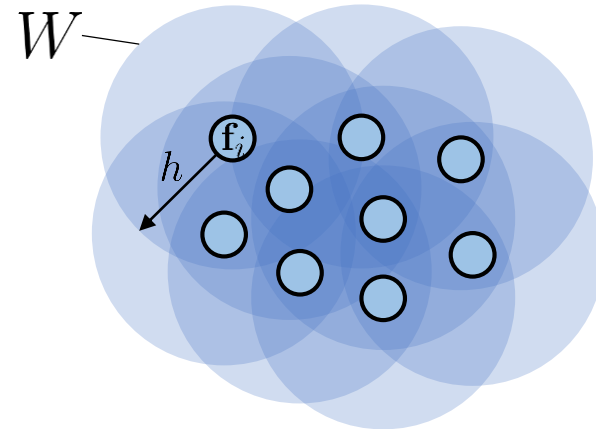
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“SPH can only be used for fluid dynamics (it’s in the name!)”

What is SPH?

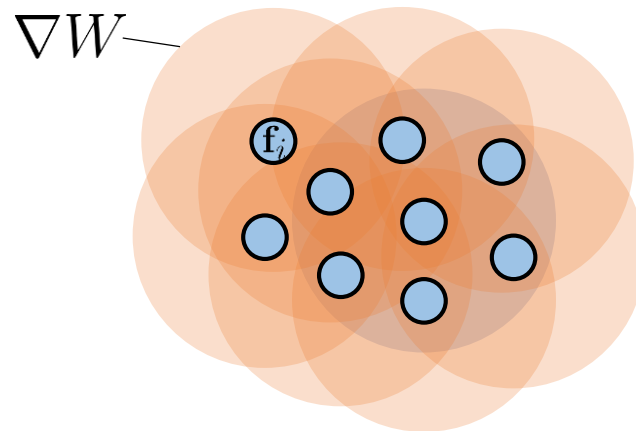
“A mesh-free method for the discretization of spatial field quantities and spatial differential operators”

$$\mathbf{f}(\mathbf{x})$$



$$\sum_i V_i \mathbf{f}_i W(\|\mathbf{x} - \mathbf{x}_i\|, h)$$

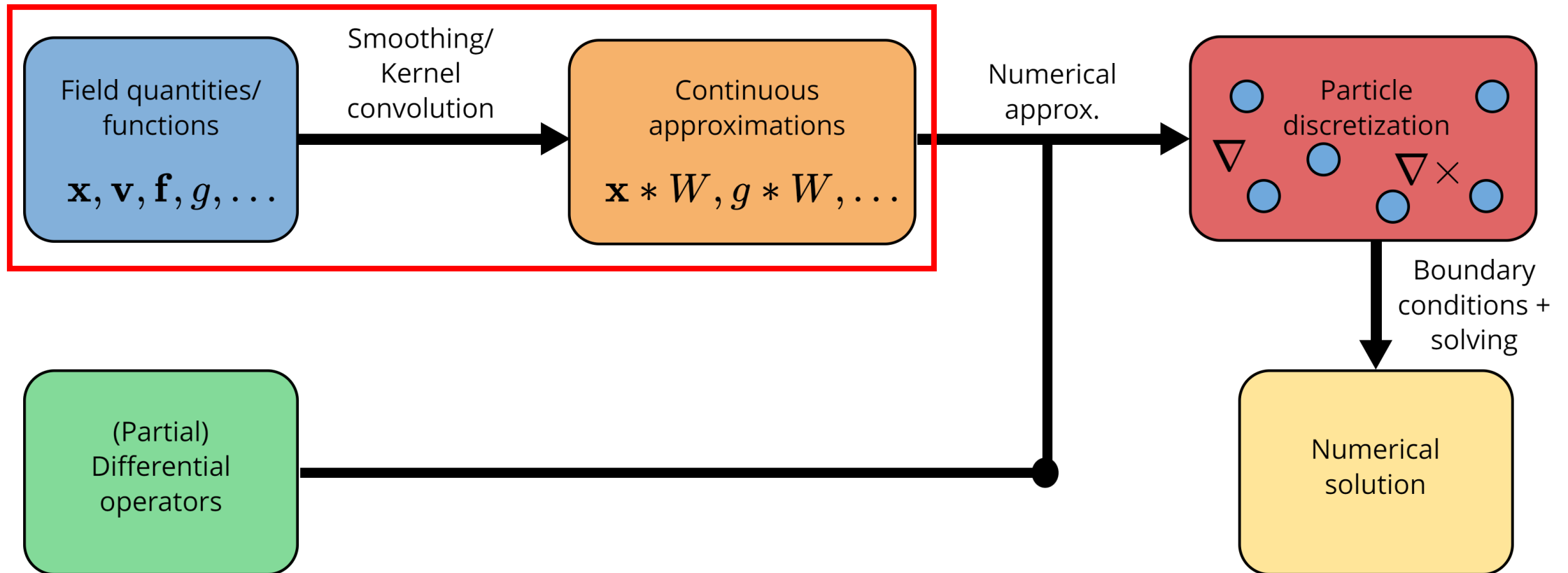
$$\nabla \mathbf{f}$$



$$\sum_i V_i \mathbf{f}_i \nabla W(\|\mathbf{x} - \mathbf{x}_i\|, h)$$

SPH Discretization Pipeline

$$\text{PDE: } \nabla \times \mathbf{f} + \nabla g = \mathbf{h}(f, g)$$



Continuous Approximation

- The dirac- δ identity $\delta(\mathbf{r}) = \begin{cases} \infty & \text{if } \mathbf{r} = 0 \\ 0 & \text{otherwise} \end{cases}$, $\int_{\mathbb{R}^d} \delta(\mathbf{r}) dv = 1$

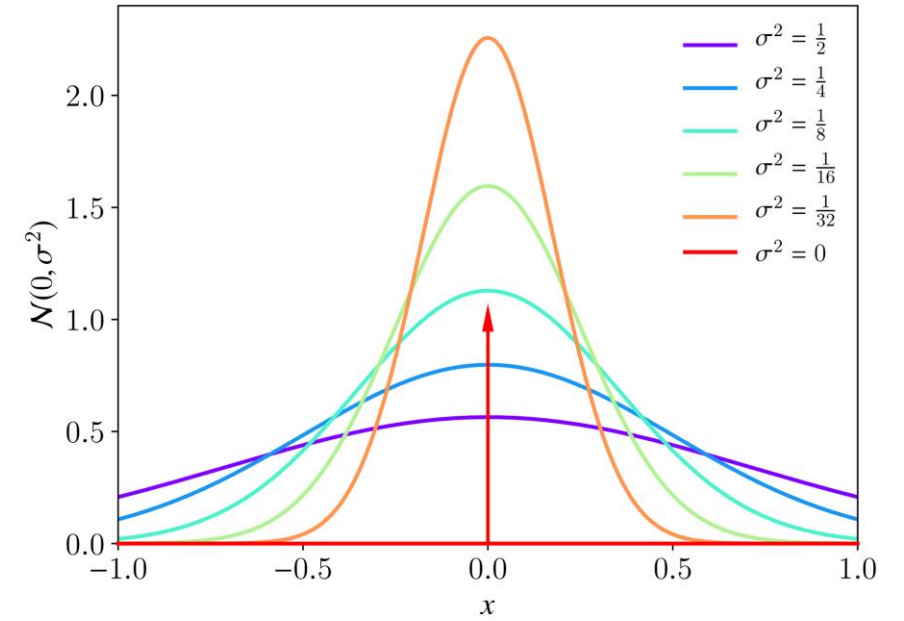
$$A(\mathbf{x}) = (A * \delta)(\mathbf{x}) = \int_{\mathbb{R}^d} A(\mathbf{x}') \delta(\mathbf{x} - \mathbf{x}') dv$$

- Approximation with Gaussian kernel $\mathcal{N}(\mathbf{x}; 0, \sigma^2)$

- Gaussian zero-variance limit $\lim_{\sigma \rightarrow 0} \mathcal{N}(\mathbf{x}; 0, \sigma^2) = \delta(\mathbf{x})$

- Replace Gaussian with kernel $W(\mathbf{x}, h) \approx \mathcal{N}(\mathbf{x}; 0, h)$

- Why is the convolution necessary? $\nabla A? \xrightarrow{\text{(shift onto kernel)}} \int_{\mathbb{R}^d} A(\mathbf{x}') \nabla W(\mathbf{x} - \mathbf{x}', h) d\mathbf{x}'$



Continuous Approximation – Kernel Choice

Essential Properties

$$\lim_{h' \rightarrow 0} W(\mathbf{r}, h') = \delta(\mathbf{r}) \quad (\text{Dirac delta condition})$$

$$\int_{\mathbb{R}^d} W(\mathbf{r}', h) = 1 \quad (\text{normalization condition})$$

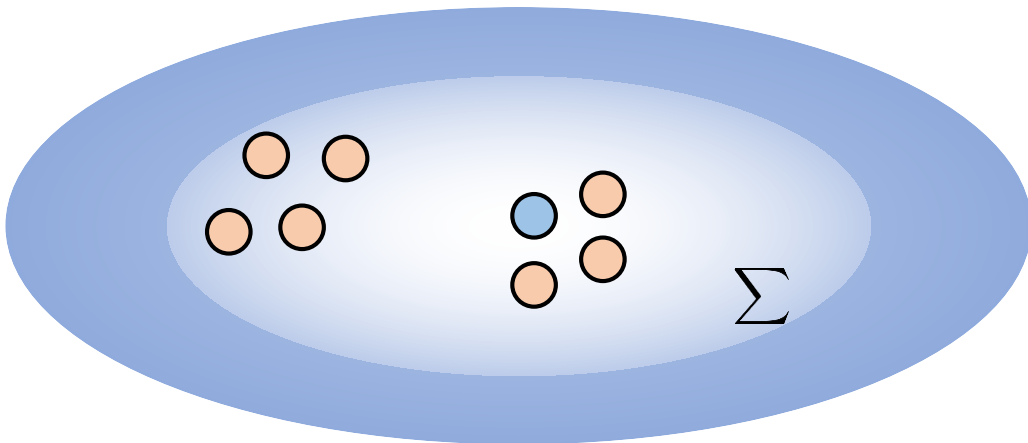
Useful Properties (optional)

$$W(\mathbf{r}, h) \geq 0 \quad (\text{positivity condition})$$

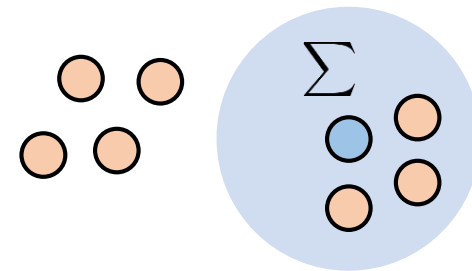
$$W(\mathbf{r}, h) = W(-\mathbf{r}, h) \quad (\text{symmetry condition})$$

$$W(\mathbf{r}, h) = 0 \text{ for } \|\mathbf{r}\| \geq \bar{h} \quad (\text{compact support cond.})$$

Global vs Compact Support



vs



Continuous Approximation – Consistency

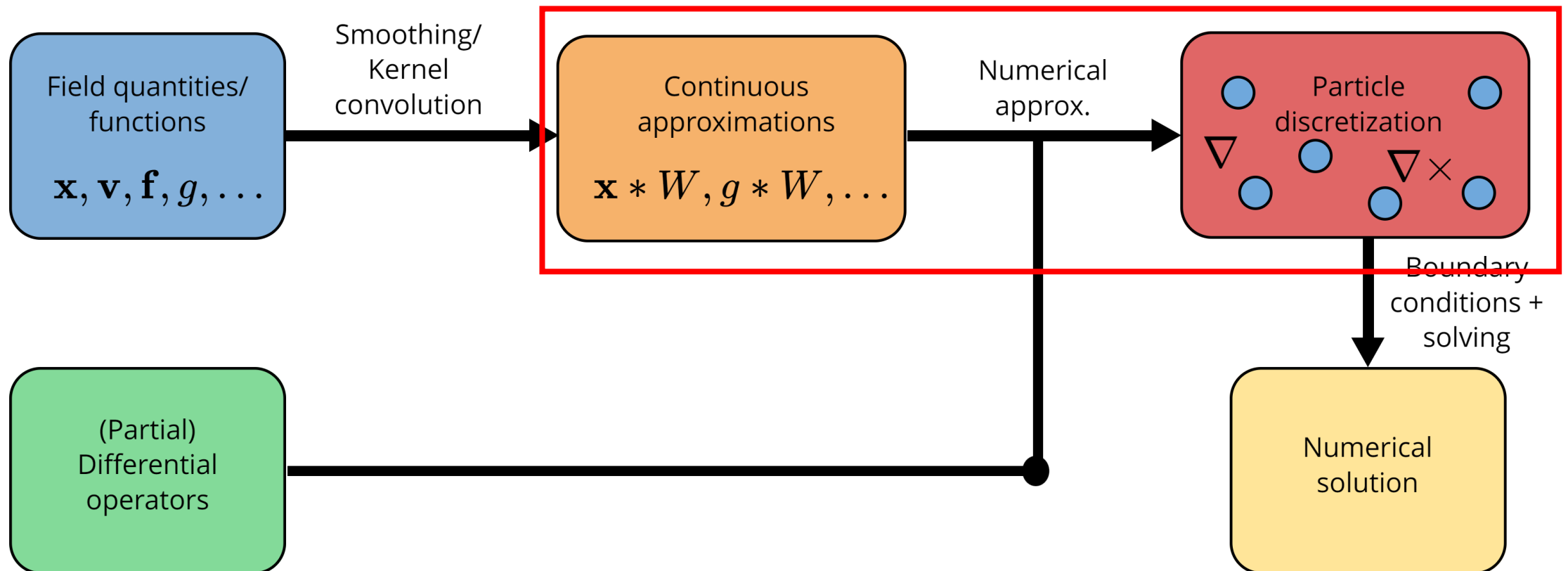
- How accurate is the continuous approximation?

$$(A * W)(\mathbf{x}) = A(\mathbf{x}) \underbrace{\int W(\mathbf{x} - \mathbf{x}', h) dv'}_{=1 \text{ if kernel normalized}} + \nabla A|_{\mathbf{x}} \cdot \underbrace{\int (\mathbf{x} - \mathbf{x}') W(\mathbf{x} - \mathbf{x}', h) dv'}_{=0 \text{ if kernel symmetric}} + \mathcal{O}(\|\mathbf{x} - \mathbf{x}'\|^2)$$

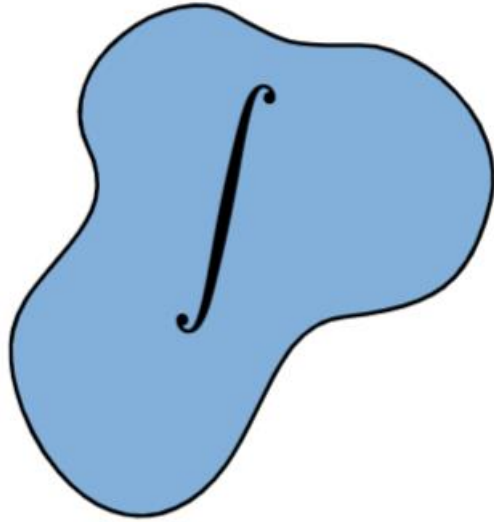
- Normalized, symmetric kernel is 1st-order consistent
- Specialized kernels with higher-order consistency can be constructed (see [LL10])

SPH Discretization Pipeline

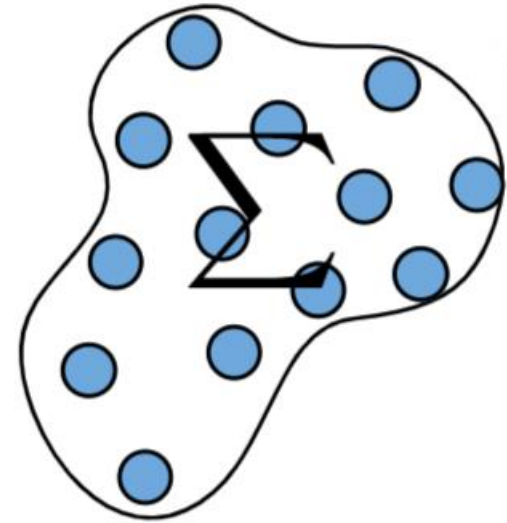
PDE: $\nabla \times \mathbf{f} + \nabla g = \mathbf{h}(f, g)$



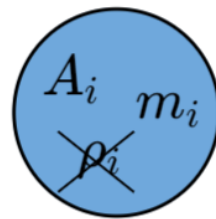
Particle Discretization



Monte-Carlo-like integration



$$(A * W)(\mathbf{x}) = \int \frac{A(\mathbf{x}')}{\rho(\mathbf{x}')} W(\mathbf{x} - \mathbf{x}', h) \underbrace{\rho(\mathbf{x}') dv'}_{dm'} \approx \sum_i A_i \frac{m_i}{\rho_i} W(\mathbf{x} - \mathbf{x}_i, h) := \langle A \rangle$$



$$\Rightarrow \rho(\mathbf{x}) \approx \sum_j m_j W(\mathbf{x} - \mathbf{x}_j, h)$$

Particle Discretization – Consistency

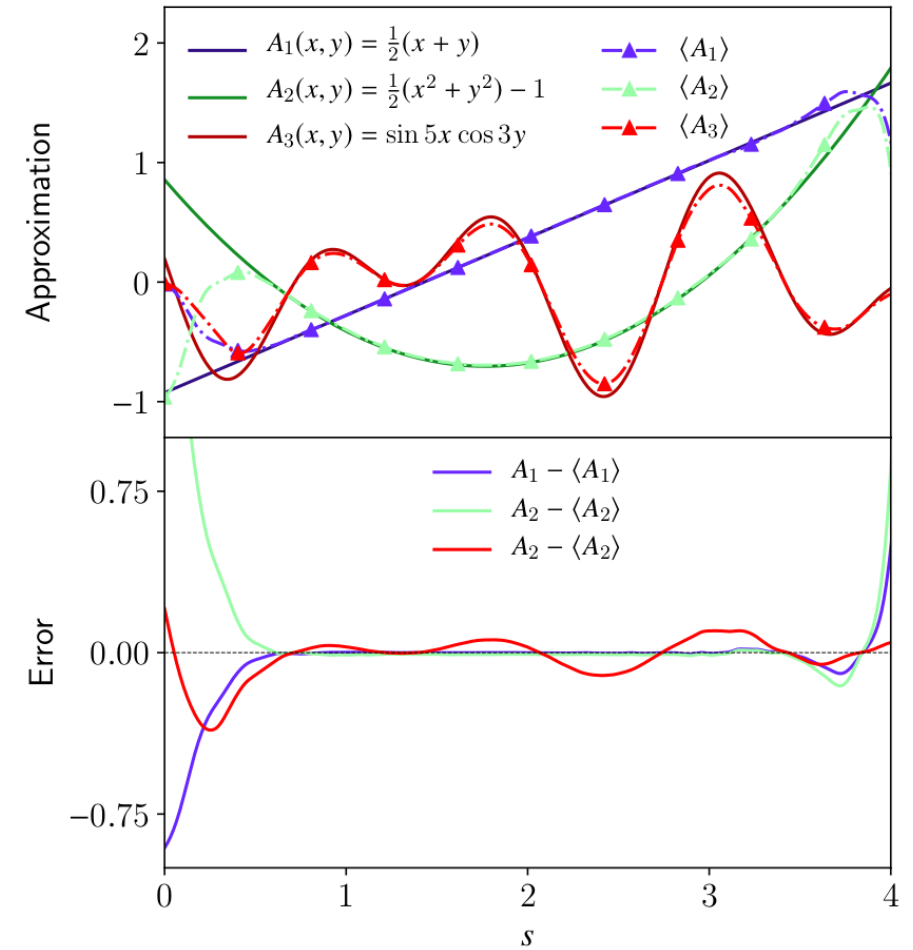
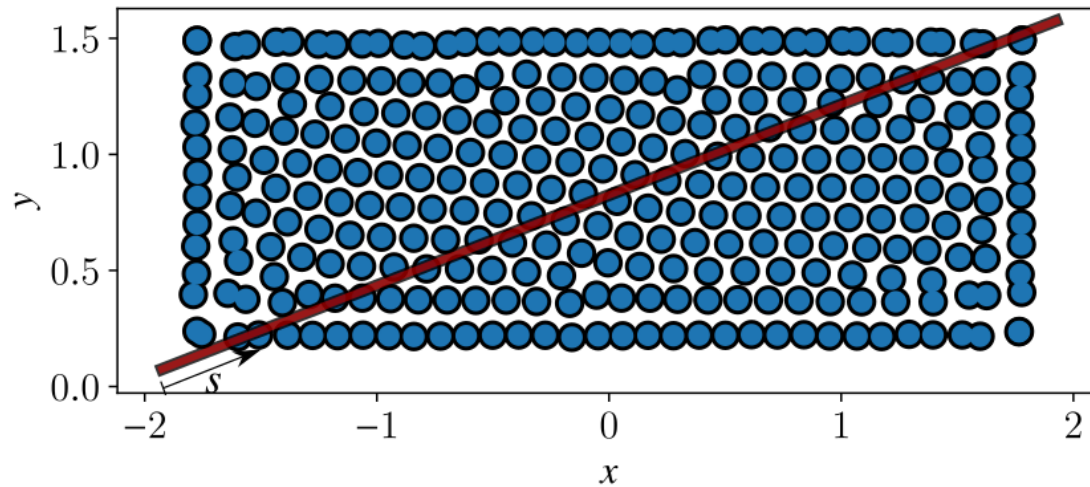
- How accurate is the discrete approximation?

$$A(\mathbf{x}_i) \approx \sum_j A_j \frac{m_j}{\rho_j} W_{ij} = A_i \underbrace{\sum_j \frac{m_j}{\rho_j} W_{ij}}_{=1 \text{ for } 0^{\text{th}}\text{-order consistency}} + \nabla A|_{\mathbf{x}_i} \cdot \underbrace{\sum_j \frac{m_j}{\rho_j} (\mathbf{x}_j - \mathbf{x}_i) W_{ij}}_{=0 \text{ for } 1^{\text{st}}\text{-order consistency}} + \mathcal{O}(\|\mathbf{x}_j - \mathbf{x}_i\|^2)$$

- Highly dependent on particle ordering
- Is this a problem?
 - Not really! Approximation quality usually still sufficient
 - Alternatively: order recovery
- Also: polynomial error is only half the truth
 - Does not take important physical conservational properties into account

Particle Discretization – Example

- **Setting:**
 - Rectangular domain discretized into particles
 - Test function sampled on particles
 - Discretization quality measured along red line
- Results good despite lack of consistency order



Operator Discretization

- Direct discretization
 - Operator “shifts” to kernel function
 - Simple
 - Gradient can be reused in implementations
- Direct form might lead to unstable simulations!
- Improved variants:
 - Difference formula
 - Symmetric formula
 - ...

$$\nabla A_i \approx \sum_j A_j \frac{m_j}{\rho_j} \nabla W_{ij}$$

$$\nabla \mathbf{A}_i \approx \sum_j \mathbf{A}_j \frac{m_j}{\rho_j} \otimes \nabla W_{ij}$$

$$\nabla \cdot \mathbf{A}_i \approx \sum_j \mathbf{A}_j \frac{m_j}{\rho_j} \cdot \nabla W_{ij}$$

$$\nabla \times \mathbf{A}_i \approx - \sum_j \mathbf{A}_j \frac{m_j}{\rho_j} \times \nabla W_{ij}$$

Symmetric Formula

- Derived from discrete Lagrangian in hydrodynamic systems:

$$\nabla A_i \approx \rho_i \left(\frac{A_i}{\rho_i^2} \langle \nabla \rho \rangle + \langle \nabla \left(\frac{A_i}{\rho_i} \right) \rangle \right) = \rho_i \sum_j m_j \left(\frac{A_i}{\rho_i^2} + \frac{A_j}{\rho_j^2} \right) \nabla_i W_{ij}$$

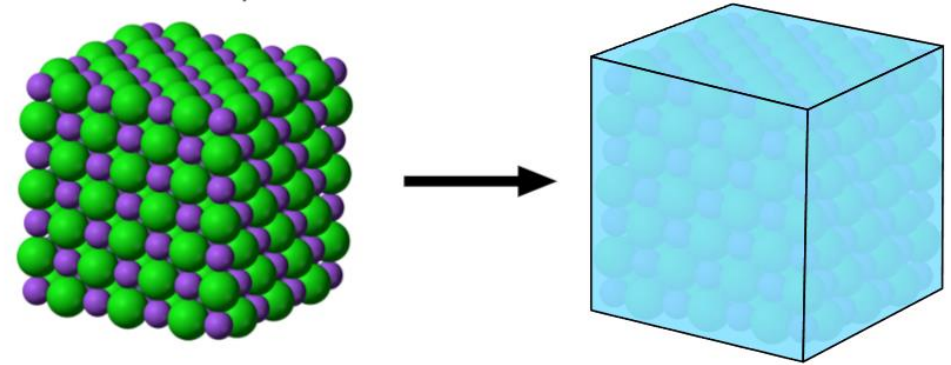
- Approx. not even 0th-order consistent. **BUT:**
 - Momentum conserving
 - Error is guided by particle ordering (non-linear properties, symmetries, conservation, etc.)
 - => Leads to more stable simulations (at least in the momentum-conservation setting)
- Consequence:

Polynomial error analysis is only half the truth

- More information: [Price2012], Smoothed Particle Hydrodynamics and Magnetohydrodynamics
 - Sec. 5 “Why a bad derivative leads to good derivatives: The importance of local conservation”

Continuum Mechanical Models

- How to model fluids and solids?
- Graphics applications:
 - Visual appearance
 - Macroscopic behavior
- What is Continuum Mechanics?
 - Continuously distributed mass
 - Object can be infinitely often divided
 - Models physical phenomena as PDE, e.g.,
 - Linear or (hyper-)elasticity
 - Navier-Stokes eq.
 - Continuum eq.



$$\rho \frac{D^2 \mathbf{x}}{Dt^2} = \nabla \cdot \mathbf{T} + \mathbf{f}_{\text{ext}}$$

- Linear momentum conservation
- $\frac{D(\cdot)}{Dt}$ denotes the material derivative
- Functions modelling \mathbf{T} are called constitutive relations

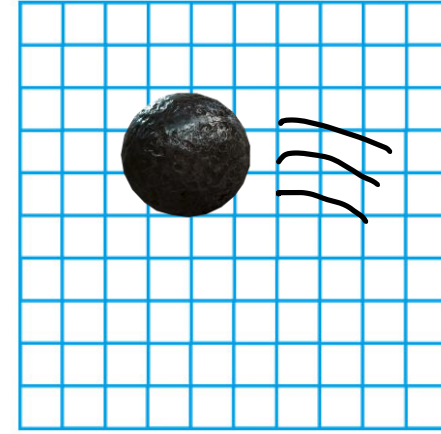
Lagrangian vs. Eulerian Point-of-View



Lagrangian Coordinates

- Identify or label material of fluid/solid
- Track material particles as they move
- Monitor changes in its properties
- Field: $A_p^L(t)$

$$\frac{DA^L}{Dt} = \frac{\partial A^L}{\partial t}$$



Eulerian Coordinates

- Identify (or label) fixed locations
- Observe fixed location (like a sensor)
- Monitor changes in properties at these
- Field: $A^E(t, \mathbf{x})$

$$\frac{DA^E}{Dt} = \frac{\partial A^E}{\partial t} + \mathbf{v} \cdot \nabla_{\mathbf{x}} A^E$$